

## Table Entries: Derivative Rules

### 1. $t$ -derivative rule

This is a course on differential equations. We should try to compute  $\mathcal{L}(f')$ . (We use the notation  $f'$  instead of  $\dot{f}$  simply because we think the dot does not sit nicely over the tall letter  $f$ .)

As usual, let  $\mathcal{L}(f)(s) = F(s)$ . Let  $f'$  be the *generalized* derivative of  $f$ . (Recall, this means jumps in  $f$  produce delta functions in  $f'$ .) The  $t$ -derivative rule is

$$\mathcal{L}(f') = sF(s) - f(0^-) \quad (1)$$

$$\mathcal{L}(f'') = s^2F(s) - sf(0^-) - f'(0^-) \quad (2)$$

$$\mathcal{L}(f^{(n)}) = s^nF(s) - s^{n-1}f(0^-) - s^{n-2}f'(0^-) + \dots + f^{(n-1)}(0^-). \quad (3)$$

**Proof:** Rule (1) is a simple consequence of the definition of Laplace transform and integration by parts.

$$\begin{aligned} \mathcal{L}(f') &= \int_{0^-}^{\infty} f'(t)e^{-st} dt & u &= e^{-st} & v' &= f'(t) \\ & & u' &= -se^{-st} & v &= f(t) \\ &= \left[ f(t)e^{-st} \right]_{0^-}^{\infty} + s \int_{0^-}^{\infty} f(t)e^{-st} dt \\ &= -f(0^-) + sF(s). \end{aligned}$$

The last equality follows from:

1. We assume  $f(t)$  has exponential order, so if  $\text{Re}(s)$  is large enough  $f(t)e^{-st}$  is 0 at  $t = \infty$ .
2. The integral in the second term is none other than the Laplace transform of  $f(t)$ .

Rule (2) follows by applying rule (1) twice.

$$\begin{aligned} \mathcal{L}(f'') &= s\mathcal{L}(f') - f'(0^-) \\ &= s(\mathcal{L}(f) - f(0^-)) - f'(0^-) \\ &= sF(s) - sf(0^-) - f'(0^-). \end{aligned}$$

Rule (3) Follows by applying rule (1)  $n$  times.

**Notes:** 1. We will call the terms  $f(0^-)$ ,  $f'(0^-)$  the 'annoying terms'. We will be happiest when our signal  $f(t)$  has rest initial conditions, so all of

the annoying terms are 0.

2. A good way to think of the  $t$ -derivative rules is

$$\begin{aligned}\mathcal{L}(f) &= F(s) \\ \mathcal{L}(f') &= sF(s) + \text{annoying terms at } 0^- \\ \mathcal{L}(f'') &= s^2F(s) + \text{annoying terms at } 0^-.\end{aligned}$$

Roughly speaking, Laplace transforms differentiation in  $t$  to multiplication by  $s$ .

3. The proof of rule (1) uses integration by parts. This is clearly valid if  $f'(t)$  is continuous at  $t = 0$ . It is also true (although we won't show this) if  $f'(t)$  is a generalized function. –See example 2 below.

**Example 1.** Let  $f(t) = e^{at}$ . We can compute  $\mathcal{L}(f')$  directly and by using rule (1).

Directly:  $f'(t) = ae^{at} \Rightarrow \mathcal{L}(f') = a/(s - a)$ .

Rule (1):  $\mathcal{L}(f) = F(s) = 1/(s - a) \Rightarrow \mathcal{L}(f') = sF(s) - f(0^-) = s/(s - a) - 1 = a/(s - a)$ .

Both methods give the same answer.

**Example 2.** Let  $u(t)$  be the unit step function, so  $\dot{u}(t) = \delta(t)$ .

Directly:  $\mathcal{L}(\dot{u}) = \mathcal{L}(\delta) = 1$ .

Rule (1):  $\mathcal{L}(\dot{u}) = s\mathcal{L}(u) - u(0^-) = s(1/s) - 0 = 1$ .

Both methods give the same answer.

**Example 3.** Let  $f(t) = t^2 + 2t + 1$ . Compute  $\mathcal{L}(f'')$  two ways.

**Solution.** Directly:  $f''(t) = 2 \Rightarrow \mathcal{L}(f'') = 2/s$ .

Using rule (3):  $\mathcal{L}(f'') = s^2F(s) - sf(0^-) - f'(0^-) = s^2(2/s^3 + 2/s^2 + 1/s) - s \cdot 1 - 2 = 2/s$ .

Both methods give the same answer.

## 2. $s$ -derivative rule

There is a certain symmetry in our formulas. If derivatives in time lead to multiplication by  $s$  then multiplication by  $t$  should lead to derivatives in  $s$ . This is true, but, as usual, there are small differences in the details of the formulas.

The  $s$ -derivative rule is

$$\mathcal{L}(tf)(s) = -F'(s) \tag{4}$$

$$\mathcal{L}(t^n f)(s) = (-1)^n F^{(n)}(s) \tag{5}$$

$$\tag{6}$$

**Proof:** Rule (4) is a simple consequence of the definition of Laplace transform.

$$\begin{aligned}
 F(s) &= \mathcal{L}(f) = \int_{0^-}^{\infty} f(t)e^{-st} dt \\
 \Rightarrow F'(s) &= \frac{d}{ds} \int_{0^-}^{\infty} f(t)e^{-st} dt \\
 &= \int_{0^-}^{\infty} -tf(t)e^{-st} dt \\
 &= -\mathcal{L}(tf(t)).
 \end{aligned}$$

Rule (5) is just rule (4) applied  $n$  times.

**Example 4.** Use the  $s$ -derivative rule to find  $\mathcal{L}(t)$ .

**Solution.** Start with  $f(t) = 1$ , then  $F(s) = 1/s$ . The  $s$ -derivative rule now says  $\mathcal{L}(t) = -F'(s) = 1/s^2$ —which we know to be the answer.

**Example 5.** Use the  $s$ -derivative rule to find  $\mathcal{L}(te^{at})$  and  $\mathcal{L}(t^n e^{at})$ .

**Solution.** Start with  $f(t) = e^{at}$ , then  $F(s) = 1/(s - a)$ . The  $s$ -derivative rule now says  $\mathcal{L}(te^{at}) = -F'(s) = 1/(s - a)^2$ .

Continuing:  $\mathcal{L}(t^2 e^{at}) = F''(s) = 2/(s - a)^3$ ,  
 $\mathcal{L}(t^3 e^{at}) = -F'''(s) = 3 \cdot 2/(s - a)^4$ ,  $\mathcal{L}(t^4 e^{at}) = F^{(4)}(s) = 4 \cdot 3 \cdot 2/(s - a)^5$ ,  
 $\mathcal{L}(t^n e^{at}) = (-1)^n F^{(n)}(s) = n!/(s - a)^{n+1}$ .

With Laplace, there is often more than one way to compute. We know  $\mathcal{L}(t^n) = n!/s^{n+1}$ . Therefore the  $s$ -shift rule also gives the above formula for  $\mathcal{L}(t^n e^{at})$ .

### 3. Repeated Quadratic Factors

Recall the table entries for repeated quadratic factors

$$\mathcal{L}\left(\frac{1}{2\omega^3}(\sin(\omega t) - \omega t \cos(\omega t))\right) = \frac{1}{(s^2 + \omega^2)^2} \quad (7)$$

$$\mathcal{L}\left(\frac{t}{2\omega} \sin(\omega t)\right) = \frac{s}{(s^2 + \omega^2)^2} \quad (8)$$

$$\mathcal{L}\left(\frac{1}{2\omega}(\sin(\omega t) + \omega t \cos(\omega t))\right) = \frac{s^2}{(s^2 + \omega^2)^2} \quad (9)$$

Previously we proved these formulas using partial fractions and factoring the denominators on the frequency side into complex linear factors. Let's prove them again using the  $s$ -derivative rule.

**Proof of (8) using the  $s$ -derivative rule.**

Let  $f(t) = \sin(\omega t)$ . We know  $F(s) = \frac{\omega}{s^2 + \omega^2}$ . The  $s$ -derivative rule implies

$$\mathcal{L}(t \sin \omega t) = -F'(s) = \frac{2\omega s}{(s^2 + \omega^2)^2}.$$

This formula is (8) with the factor of  $2\omega$  moved from one side to the other.

The other two formulas can be proved in a similar fashion. We won't give the proofs here.

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